

## PASSIVATING MULTI CRISTALLINE Si SOLAR CELLS USING SiN<sub>x</sub>:H

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### ABSTRACT:

Passivating solar cells with SiN<sub>x</sub>:H has been so far a scarcely understood effect that can only be optimized for cell production in an empirical way. In this study we determine the structural properties of SiN<sub>x</sub>:H layers with Fourier Transform Infrared (FTIR) measurements and relate these to both the deposition parameters and its passivating qualities for solar cells. Furthermore we determined the relations between the hydrogen diffusion in the SiN<sub>x</sub>:H and the structural properties of these layers.

The Si-N, Si-H and N-H bond densities for layers deposited with either nitrogen (N<sub>2</sub>), ammonia (NH<sub>3</sub>) or deuterated ammonia (ND<sub>3</sub>) and silane (SiH<sub>4</sub>) are affected by the N/Si flow ratio and the pressure  $p$  in a similar way although the differences in dissociation energy and rate cause different deposition mechanisms. Comparing the Si-N and Si-H bond densities found for NH<sub>3</sub> and ND<sub>3</sub> grown layers, we find that roughly 25% of the hydrogen in the SiN<sub>x</sub>:H layers stems from the ammonia precursor gas, while 75% stems from the silane.

We show that Si-N bond density is an important parameter governing both the bulk and surface passivation of the SiN<sub>x</sub>:H layers. In spite of the different deposition mechanisms, the same relations hold between the H-diffusion coefficient, Si-N bond density and passivating qualities of SiN<sub>x</sub>:H layers deposited with either N<sub>2</sub> or NH<sub>3</sub>. The best bulk and surface passivating layers have a relatively low hydrogen diffusion coefficient due to a high Si-N bond density. We find optimum values for bulk and surface passivation for Si-N bond densities of  $1.3 \cdot 10^{23} \text{ cm}^{-3}$ , regardless of the type of SiN<sub>x</sub>:H used and regardless of the starting wafer quality. Lower Si-N bond densities result in layers with a more open structure and this will probably lead to H<sub>2</sub> formation during annealing. These H<sub>2</sub> molecules will effuse into the ambient during firing, and do not contribute to the passivation of solar cells. Higher Si-N bond densities result in a too dense structure, prohibiting an effective diffusion of H-atoms into the bulk of the solar cells. This study further indicates that FTIR analysis gives us a quick and reliable tool to check the quality and properties of SiN<sub>x</sub>:H layers. This will allow optimization of SiN<sub>x</sub>:H deposition systems without having to make complete solar cells.

## 1 INTRODUCTION

Hydrogenated amorphous silicon nitride layers (a-SiN<sub>x</sub>:H) have become a very important part of modern silicon PV technology. They act as an anti-reflection coating and provide both bulk and surface passivation, important means for optimizing multi crystalline (mc) Si solar cells and obtaining high efficiencies [1]. The bulk, or defect passivation is achieved by driving hydrogen into the mc Si solar cells by a short thermal anneal [2,3,4]. At the same time, the surface passivation is improved by reducing the number of dangling bonds and creating a positive field effect by K<sup>+</sup> centers [5].

One of the key issues of our research is to combine excellent bulk and surface passivation on low cost silicon with easy-to-handle gasses. Relations between structural properties and the quality of bulk and surface passivation of SiN<sub>x</sub>:H need to be known for better understanding of the underlying physics. So far, passivation with SiN<sub>x</sub>:H in industry has been a phenomenon that can only be optimized for cell production in an empirical way. This study provides means to determine passivating qualities of SiN<sub>x</sub>:H independently of the deposition method and without making solar cells.

The amount of hydrogen diffusion from the SiN<sub>x</sub>:H layer into the silicon wafer largely determines the quality of bulk passivation. Numerous generic studies of hydrogen diffusion in silicon nitride films have been reported [6,7,8,9], but studies on the relation between hydrogen diffusion in silicon nitride and bulk passivation are rather scarce [10,11]. The surface passivation is commonly related to the refractive index [5,12] and not to the structural properties. Only Mäckel and Lüdemann [13] and ECN [14] performed such a study. In recent publications, Weeber et al. [14,15] presented a systematic investigation on the bulk and surface passivating properties, the hydrogen loss mechanism in the SiN<sub>x</sub>:H layer, and related these to the structural properties of the SiN<sub>x</sub>:H deposited with SiH<sub>4</sub> and N<sub>2</sub>. This first study is now extended to the use of NH<sub>3</sub> and deuterated ammonia (ND<sub>3</sub>). The use of ND<sub>3</sub> sheds more light on the plasma chemistry and the deposition mechanisms of the SiN<sub>x</sub>:H layer, while the use of both N<sub>2</sub> and NH<sub>3</sub> nitrides on solar cells reveals the structural properties of the SiN<sub>x</sub>:H

layers that govern its inherent passivating qualities independent of the deposition mechanisms.

## 2 EXPERIMENTAL

### 2.1 Approach

The physics behind the passivating properties of  $\text{SiN}_x\text{:H}$  layers is investigated in four steps:

- Deposition of layers on double polished Cz Si wafers for FTIR analysis to obtain bond densities, and relating those to the deposition parameters;
- Application of the examined layers to solar cells; determination of bulk passivating quality by measuring  $V_{oc}$  and the Internal Quantum Efficiency at 1000 nm ( $\text{IQE}_{(1000\text{nm})}$ ), and relating these to the structural properties;
- Application of the examined layers to FZ Si wafers; determination of surface passivation by measuring the effective lifetime ( $\tau_{\text{eff}}$ ) using the Quasi Steady State Photo Conductance method [16] and relating these to the structural properties;
- Annealing of  $\text{SiN}_x$  layers for increasing time periods at 800 °C to determine H-diffusivity in the  $\text{SiN}_x\text{:H}$  layer and relate that to the structural properties of  $\text{SiN}_x\text{:H}$ ;

### 2.2 Deposition of $\text{SiN}_x\text{:H}$ layers

We deposited  $\text{SiN}_x\text{:H}$  layers with three different types of nitrogen-containing precursor gases; nitrogen ( $\text{N}_2$ ), ammonia ( $\text{NH}_3$ ) and deuterated ammonia ( $\text{ND}_3$ ). This will allow us to determine whether fundamental relations exist independent of the process gases. Silane ( $\text{SiH}_4$ ) was used as silicon precursor gas. The  $\text{SiN}_x\text{:H}$  layers were deposited using the in-line microwave remote Plasma Enhanced Chemical Vapor Deposition (MW RPECVD) system at ECN at different process parameters such as pressure and gas flows. This MW RPECVD has been co-developed with Roth and Rau [17].

### 2.3 FTIR analysis

We determined the bond densities (Si-H, Si-N and N-H) within the nitride layers using FTIR spectroscopy. From the transmission spectrum of the  $\text{SiN}_x\text{:H}$  layers deposited on Si substrates the absorption spectrum ( $k$ ) of the single layer of  $\text{SiN}_x\text{:H}$  was calculated according to the method reported by Maley [18]. Subsequently, the bond densities are calculated by integrating the different absorption peaks and multiplying them with a proper proportionality constant [19,20].

### 2.4 Solar cells

The mc Si solar cells for the determination of the bulk passivation were made using neighboring wafers; the process sequence was: texturing using acidic etching, emitter formation using an infrared heated belt furnace, remote MW PECVD of  $\text{SiN}_x\text{:H}$  using the same deposition parameters as for the FTIR samples, metallization using screen printing and contact formation with an infrared heated belt furnace. IV measurements were performed with our Class A solar simulator according to the ASTM-E948 norm; the Internal Quantum Efficiency (IQE) was calculated from the spectral response and the reflectance.

## 3. EXPERIMENTAL RESULTS AND DISCUSSION

### 3.1 $\text{SiN}_x\text{:H}$ structural properties related to the deposition parameters

In previous studies it was found that the mass density of the  $\text{SiN}_x\text{:H}$  layers is related to the Si-N bond density: layers containing more Si-N bonds have a higher mass density  $\rho$ , while those with relatively more Si-Si bonds are less dense and have a more porous structure [14, 21, 22].

In figure 1 and 2 the Si-N bond density for three nitrides deposited with  $\text{NH}_3$ ,  $\text{ND}_3$  and  $\text{N}_2$  as precursor gasses are shown as a function of the N/Si gas flow ratio, and for different deposition pressures. For all nitrides the Si-N bond densities increase for increasing N/Si flow ratio, and are higher for lower deposition pressures. The refraction index  $n$  increases with decreasing Si-N bond density and decreasing N/Si flow ratio, which is shown qualitatively by the arrow in the graphs.

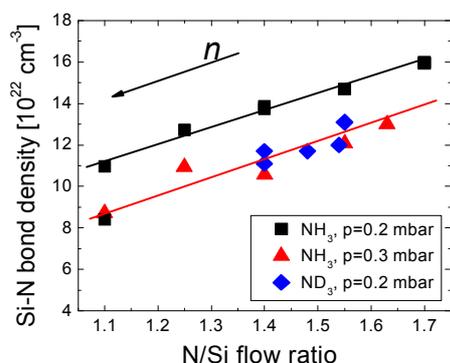


Figure 1: Si-N bond density versus the N/Si flow ratio for nitrides deposited with  $\text{NH}_3$  and  $\text{ND}_3$  at different pressures.

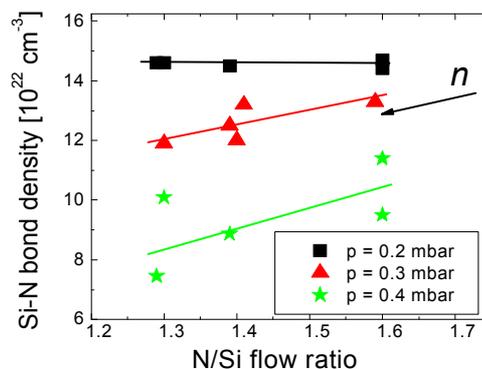


Figure 2: Si-N bond density vs N/Si flow ratio for nitrides deposited with  $\text{N}_2$  at different pressures.

In figure 3 and 4 the Si-H and N-H bond densities of the same  $\text{SiN}_x\text{H}$  layers are shown. For the Si-H and N-H bond densities of the  $\text{NH}_3$  and  $\text{ND}_3$  nitrides (figure 3) no pressure dependence is found; the Si-H bond density decreases, while the N-H bond density increases with increasing gas flow ratio. The Si-H and N-H bond densities of the  $\text{N}_2$  nitrides however, do not significantly depend on the N/Si flow ratio. For these bond densities we see the pressure dependence as a main effect repeated. The points in figure 4 are averages over all values at a certain pressure.

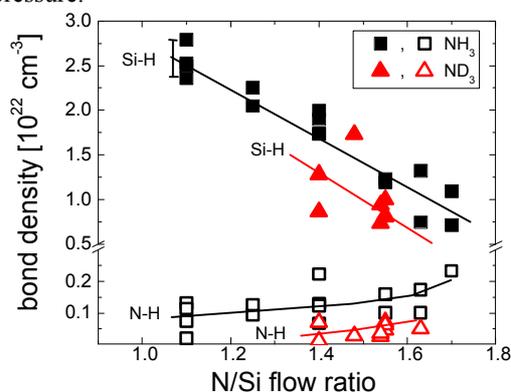


Figure 3: Si-H and N-H bond density versus the N/Si flow ratio for the  $\text{NH}_3$  and  $\text{ND}_3$  nitrides.

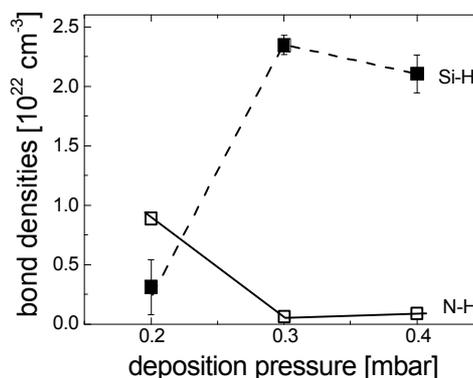


Figure 4: Si-H and N-H bond density versus the deposition pressure for  $\text{N}_2$  nitrides.

The bond densities of the nitrides deposited with  $\text{NH}_3$  (and  $\text{ND}_3$ ) are stronger dependent on the flow ratio than those of nitrides deposited with  $\text{N}_2$ , especially at lower pressures ( $p = 0.2$  mbar); the structural properties of  $\text{N}_2$  layers depend mainly on the pressure. The increase in Si-N bond density (that we see for all nitrides) with increasing N/Si flow ratio is easily understood. At higher N/Si flow ratios relatively more reactive N-containing species are available; this will favor the formation of Si-N and N-H bonds over that of Si-Si and Si-H. The decrease in Si-H bonds with increasing N/Si flow ratio is only seen for the  $\text{NH}_3$  and  $\text{ND}_3$  nitrides. Mäckel and Lüdemann [13] found that the Si-N bond density increases (and the Si-H bond density decreases) with increasing N/Si atomic ratio in the layer. This corresponds to our results, since an increasing N/Si flow ratio will increase the N/Si atomic ratio in the layer.

The decrease of Si-N bond density for increasing pressures, also seen in all nitrides, is due to properties of our remote MW PECVD. In our system, the nitrogen containing gas is fed in near the plasma source, while the silane gas is fed in just above the samples. At the microwave source nitrogen (or ammonia) gas is dissociated by electron impact; while the silane in turn is dissociated by either electron impact or by interactions with N- or H- (in case of ammonia) radicals. Besides this, the dissociation of silane can be enhanced or even caused by the high temperatures in the deposition chamber. Higher pressures will confine the nitrogen plasma closer to the plasma source, and less reactive N-containing species will reach the substrate. This will cause the formation of more Si-rich  $\text{SiN}_x\text{H}$  layers that are less dense. A more detailed description of the plasma chemistry can be found in other

publications [23,24].

The different flow and pressure dependence for both types of nitrides is caused by the difference in dissociation energy for N-N (9.81 eV) and H-NH<sub>2</sub> (4.65 eV) [25]. At low pressures, the degree of N<sub>2</sub> depletion will saturate for a certain microwave power; rendering the plasma independent on further increasing the N<sub>2</sub>/SiH<sub>4</sub> flow ratio. Due to the lower ionization energy this maximum is not reached for NH<sub>3</sub>. Within our processing parameters and the depletion of the NH<sub>3</sub> is close to 100% [17].

The lower Si-N bond density of the deuterated nitride, when compared to the nitride deposited with NH<sub>3</sub>, is due to the difference in dissociation rate of NH<sub>3</sub> and ND<sub>3</sub>. This dissociation rate depends on the vibration frequency of the N-D or N-H in the molecule, which is lower for N-D due to the larger mass. The Si-H bond density in the layers deposited with ND<sub>3</sub> is 75% of that in the corresponding layer deposited with NH<sub>3</sub>, while the N-H bond density in layers made with ND<sub>3</sub> is hardly 50%. Since (in NH<sub>3</sub> and ND<sub>3</sub> nitrides) the majority of hydrogen is bonded to Si-atoms (see figure 3), this means that at least 25% of the hydrogen in the SiN<sub>x</sub>:H layers stems from ammonia, while roughly 75% stems from silane.

In figure 5 the initial hydrogen fraction of all layers (deposited with N<sub>2</sub>, NH<sub>3</sub> or ND<sub>3</sub>) is shown against the Si-N bond density. We define the hydrogen fraction as:

$$\text{fraction H} = \frac{[\text{Si-H}] + [\text{N-H}]}{[\text{Si-H}] + [\text{N-H}] + [\text{Si-N}]} \quad (1)$$

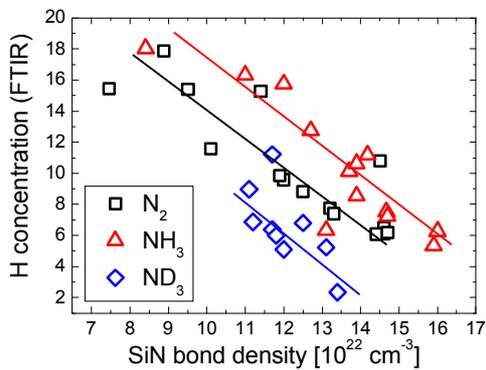


Figure 5: Total hydrogen concentration (calculated from bond densities) for the three different nitrides as a function of the Si-N bond density.

Although there is a lot of variation, the general trend is clear for all layers: with increasing Si-N bond density (and thus layer density) the total H-content decreases. For most SiN<sub>x</sub>:H layers deposited with NH<sub>3</sub>, more hydrogen is incorporated at the same Si-N bond densities than in layers deposited with N<sub>2</sub> or ND<sub>3</sub>. As was found above, in layers deposited with ND<sub>3</sub> at least 25% of the hydrogen is replaced by deuterium. But, layers deposited with ND<sub>3</sub> also become less dense (lower Si-N bond density) using the same process parameters when compared with layers deposited with NH<sub>3</sub> due to the difference in dissociation rate, as was shown in figure 1. Both effects cause the H-content in ND<sub>3</sub> layers to be only 50% of the content in NH<sub>3</sub> layers at the same Si-N bond densities.

When N<sub>2</sub> is used instead of NH<sub>3</sub>, the total H content decreases by roughly 25%. In this deposition, all hydrogen will stem from the silane; the 25% lower H-content is in agreement with the 25% difference in H-content between layers deposited with NH<sub>3</sub> and ND<sub>3</sub>. This could mean that although the plasma chemistry is different, efficiency of hydrogen incorporation is approximately the same.

### 3.2 Bulk and surface passivation as a function of Si-N bond density

In previous publications we show that the Si-N bond density is a key parameter for bulk passivation [14,15,26]. Figure 6 shows the V<sub>oc</sub> of cells made with SiN<sub>x</sub>:H layers deposited with N<sub>2</sub> or NH<sub>3</sub>. Wafers of average quality (A) and wafers of higher quality (B) were used. All values of V<sub>oc</sub> are scaled to the maximum V<sub>oc</sub> for each experimental run to be able to compare all the data.

From the figure it can be seen that:

- The V<sub>oc</sub> for cells with SiN<sub>x</sub>:H layers deposited with N<sub>2</sub> or NH<sub>3</sub> have the same dependence of the Si-N bond density.
- The maximum V<sub>oc</sub> for both material qualities is found at the same Si-N bond density, viz. 1.3\*10<sup>23</sup> cm<sup>-3</sup>. The difference is that for the better quality higher V<sub>oc</sub> values are found for a broader range of Si-N bond densities.

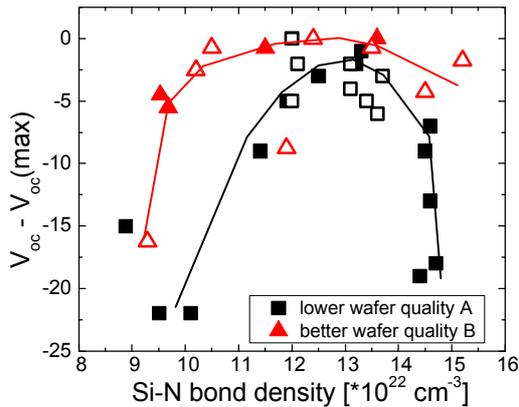


Figure 6:  $V_{oc} - V_{oc(max)}$  of mc-Si solar cells versus the Si-N bond density. The squares show the  $V_{oc}$  for cells made of average quality (A), the triangles show  $V_{oc}$  for cells made of good wafer quality (B). Closed symbols:  $N_2$  nitrides; open symbols:  $NH_3$  nitrides. The lines are guides to the eye.

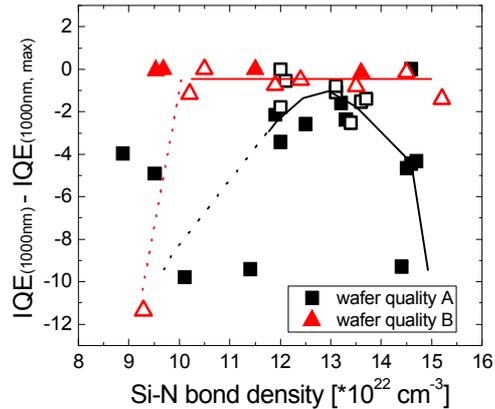


Figure 7: IQE at 1000 nm for the different solar cells. Again the difference between the different wafer qualities is clearly visible. Closed symbols:  $N_2$  nitrides; open symbols:  $NH_3$  nitrides. The lines are guides to the eye.

We were not able to make solar cells with  $[Si-N] < 9 \cdot 10^{22} \text{ cm}^{-3}$  or  $[Si-N] > 1.5 \cdot 10^{23} \text{ cm}^{-3}$ , because these layers became too inhomogeneous at larger areas.

$V_{oc}$  is determined by both the bulk and surface passivation. To investigate the dependence of the Si-N bond density on bulk passivation the IQE at 1000 nm is shown in figure 7. When the wafer quality is lower (squares in graph 6 and 7), the  $V_{oc}$  follows the IQE at 1000 nm indicating  $V_{oc}$  is mainly influenced by the bulk properties of the solar cell. In the case of better wafer quality (triangles) however, the IQE remains constant down to very low Si-N bond densities ( $1 \cdot 10^{23} \text{ cm}^{-3}$ ). In this case, the bulk properties remain constant and smaller changes in  $V_{oc}$  (<5 mV) are caused by differences in surface passivation; the better surface passivation for Si-N bond densities around  $1 \cdot 10^{23} \text{ cm}^{-3}$  will be confirmed below. At the lowest bond densities, the drop in IQE reflects a drastic reduction in bulk passivation; this results in an additional loss in  $V_{oc}$  of about 10 mV.

The amount of surface passivation can be determined by measuring the effective lifetime of charge carriers in FZ wafers coated with  $SiN_x:H$  layers. In figure 8 we show the effective lifetimes for  $N_2$  (closed symbols), and some  $NH_3$  (open symbols) grown nitrides with the different Si-N bond densities, before and after a very long anneal (60 min at 800 °C) for the  $N_2$  nitrides, and a shorter anneal for the  $NH_3$  nitrides.

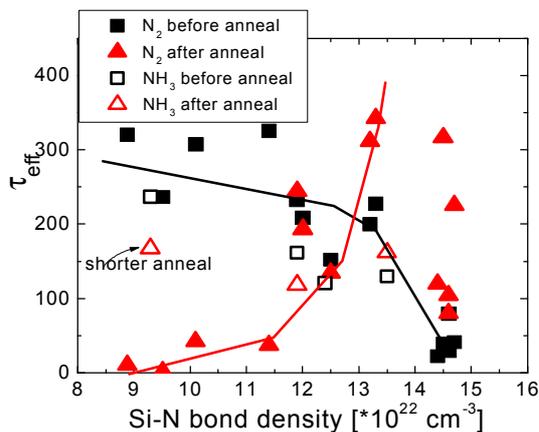


Figure 8: Effective lifetimes of charge carriers in FZ wafers with  $SiN_x:H$  layers before and after anneal (60 mins, 800 °C).  $N_2$  nitrides: closed symbols,  $NH_3$  nitrides: open symbols.

The figure shows that good initial surface passivation (large  $\tau_{eff}$ ) can be obtained for layers with low Si-N bond densities, but that this surface passivation is not thermally stable after a 60 minute anneal (although the much shorter firing of solar cells does not influence the  $\tau_{eff}$  this much, see also arrow in figure 8 [23,27]). At high Si-N bond densities, the initial  $\tau_{eff}$  is low, but after a long anneal the values improve drastically. Around  $[Si-N] \sim 1.3 \cdot 10^{23} \text{ cm}^{-3}$ ,  $\tau_{eff}$  remains constant at high values. At these bond densities,  $V_{oc}$  for wafers of good quality reaches its maximum (see figure 6, wafer quality B) that we contribute to the good and stable surface passivation. At higher Si-N bond densities there is a large variation in lifetime after annealing. This means that the best and most stable processing is found at a Si-N bond density of  $1.3 \cdot 10^{23} \text{ cm}^{-3}$ , the same density as for best bulk passivation.

Combining the results for the bulk and surface passivation, we find optimum values for Si-N bond densities of  $1.3 \cdot 10^{23} \text{ cm}^{-3}$  for both, regardless of the type of  $\text{SiN}_x\text{:H}$  used and regardless of the wafer quality.

### 3.3 Hydrogen diffusion in $\text{SiN}_x\text{:H}$ layers.

In earlier publications we reported that the best surface and bulk passivation is found for denser layers, with a higher Si-N bond density ( $1.3 \cdot 10^{23} \text{ cm}^{-3}$ ), regardless of the initial amount of hydrogen. In these studies it was postulated that low Si-N bond densities result in more open structures that facilitate the formation of  $\text{H}_2$  molecules that effuse into the ambient during anneal. This effused hydrogen will not contribute to the bulk passivation. In too dense layers ( $[\text{Si-N}] > 1.3 \cdot 10^{23} \text{ cm}^{-3}$ ) the H-diffusion into the bulk of the solar cells during short anneals will be too small for good passivation [14,15,26]. In a recent publication, Dekkers et al [11] found that the mass density of the  $\text{SiN}_x\text{:H}$  layer governs the diffusion of hydrogen in the layer, denser layers prevent out-diffusion of molecular hydrogen. This publication confirms our earlier findings.

Our previous study on the H-diffusion in  $\text{SiN}_x\text{:H}$  and the related passivating properties was performed on  $\text{SiN}_x\text{:H}$  layers deposited with  $\text{N}_2$  and  $\text{SiH}_4$ . This experiment is now extended with layers deposited with  $\text{NH}_3$  and  $\text{SiH}_4$ . The diffusion of H in the  $\text{NH}_3$  nitrides is monitored by measuring the Si-H and N-H bond densities as function of the anneal time at  $800^\circ\text{C}$ .

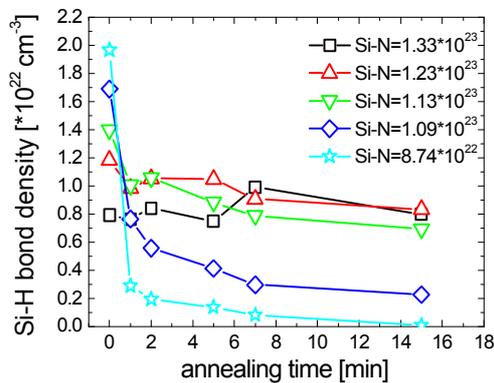


Figure 9: Change in Si-H bond density during anneal for different Si-N bond densities. All layers are deposited with  $\text{NH}_3$  and  $\text{SiH}_4$ .

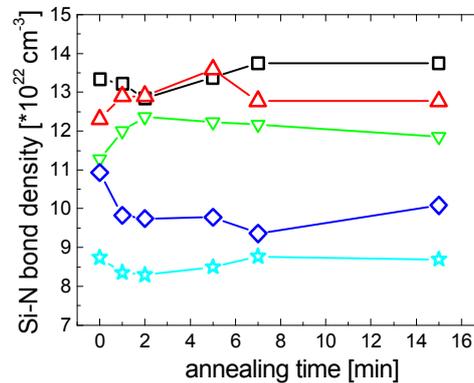


Figure 10: Change in Si-N bond density during anneal. All layers are deposited with  $\text{NH}_3$  and  $\text{SiH}_4$ . The symbols correspond to those in figure 9.

In figure 9 the change in Si-H bond density upon annealing at  $800^\circ\text{C}$  is shown for layers with different Si-N bond densities, deposited with  $\text{SiH}_4$  and  $\text{NH}_3$ . The N-H bond density shows similar behavior as a function of time. It is clear that layers with a high initial H-concentration (low Si-N bond density) lose most hydrogen during anneal. In figure 10 the change in Si-N bond density during the same anneal times is shown. Although some changes may be seen during the first anneal periods, the Si-N bond density remains approximately constant during anneal.

Similar to the  $\text{N}_2$  nitrides [14,15], also in  $\text{NH}_3$  nitrides the H-diffusion shows a strong time-dispersive (time dependent diffusion coefficient) behavior. This effect is well known for H diffusion in a-Si:H and can be described by [28]:

$$C_H = C_{H,0} \exp(-\pi^2 D_H t / L^2) \quad (2)$$

$$D_H(t) = D_H(0) (\omega t)^{-\alpha} \quad (3)$$

$\omega$  is the H attempt-to-diffuse frequency and  $\alpha$  is the temperature dependent dispersion parameter ( $0 < \alpha < 1$ ). More details about this dispersive character of H diffusion in  $\text{SiN}_x\text{:H}$  can be found in [24].

The initial diffusion coefficient  $D_1$  of the hydrogen strongly depends on the Si-N bond density in the layers. In figure 11, these coefficients for Si-H for both  $\text{N}_2$  and  $\text{NH}_3$  nitrides are shown. The coefficients for N-H bond densities are similar, but harder to measure since the peak is much smaller and the error much larger. Although the hydrogen concentrations are quite different for  $\text{N}_2$  and  $\text{NH}_3$  nitrides (see figure 5), the diffusion coefficients are the same for both materials, and only depend on the Si-N bond densities.

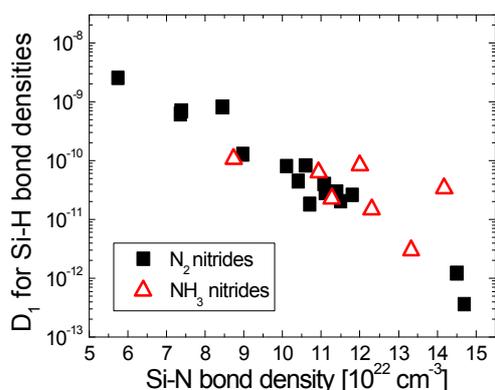


Figure 11: Initial Si-H diffusion coefficient  $D_1$  as a function of Si-N bond density.

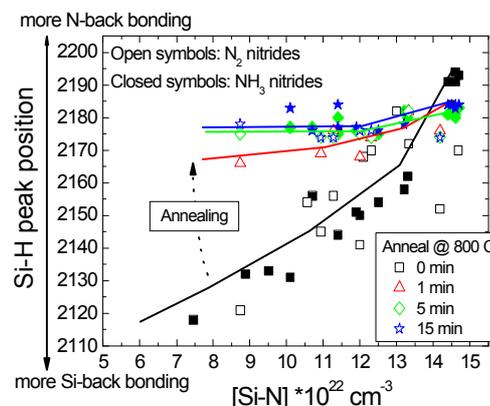


Figure 12: Shift in Si-H peak position as a function of annealing time.

A similar agreement between  $N_2$  and  $NH_3$  nitrides is also found for the shift in the Si-H peak position, shown in figure 12. The Si-H peak position is shown for both  $N_2$  and  $NH_3$  nitrides before and after annealing at several times. The Si-H peak position depends only on the Si-N bond density and the anneal times. The Si-H stretch vibration mode in  $SiN_x:H$  can occur at various frequencies between 2000 and 2300  $cm^{-1}$ , depending on the back bonding of the concerning Si atom. If the back bonds are mainly Si atoms the vibration frequency is low (2000  $cm^{-1}$  for H-Si-Si<sub>3</sub>), in the case of mainly N atoms the vibration frequency is high (2220  $cm^{-1}$  for H-Si-N<sub>3</sub>) [21,22]. As expected, before annealing ( $t = 0$ ; square symbols) we see that if the Si-N bond density is low, thus more Si in the layer, the Si-H back bonding is mainly silicon. When the Si-N bond density is high (more N in the layer), the Si-H back bonding is mainly nitrogen.

In figure 12 we see that after annealing the Si-H bonds with Si-rich back bonding disappear from all nitrides, only the Si-H bonds with N-rich back bonding remain. Indeed, figure 11 shows that  $SiN_x:H$  layers with low Si-N bond densities and thus more Si-rich back bonding have the highest Si-H diffusion coefficient. This stems from the more open and porous structure of the Si-rich layers. The strongest and most stable bonds are those around 2180 (H-Si-HN<sub>2</sub>).

Although the  $N_2$  and  $NH_3$  nitrides are formed with different deposition mechanics and the initial hydrogen concentrations are not the same, the hydrogen diffusion and corresponding structural changes are the same for layers with the same Si-N bond density. In paragraph 3.2 we already showed that the bulk passivating properties of the two layers depend in the same way on the Si-N bond density; this fact can now be explained by the similar H diffusion from the  $SiN_x:H$  layer into the mc-silicon due to the same Si-N bond density in the layers.

#### 4. CONCLUSIONS

We have shown that for  $SiN_x:H$  layers deposited with our remote MW PECVD system the N/Si flow ratio and the deposition pressure are important processing parameters. The Si-N, Si-H and N-H bond densities of layers deposited with  $N_2$ ,  $NH_3$  or  $ND_3$  are related to N/Si flow ratio and the pressure  $p$  in a similar way, but the difference in dissociation energy and rate causes differences in the deposition mechanisms. Comparing the Si-N and Si-H bond densities found for  $NH_3$  and  $ND_3$  grown layers, we find that roughly 25% of the hydrogen in the  $SiN_x:H$  layers stems from the ammonia precursor gas, while 75% stems from the silane. This fact is confirmed by the 25% difference in H-content in  $N_2$  and  $NH_3$  layers.

Even though the  $N_2$  and  $NH_3$  nitrides are formed with different deposition mechanics and the initial hydrogen concentrations are not the same, the hydrogen diffusion and thus bulk passivation are the same for layers with the same Si-N bond density. Layers with low Si-N bond densities are porous with large hydrogen diffusion coefficients. During anneal the hydrogen will effuse into the ambient in molecular form and will not contribute to the passivation. Dense layers on the other hand, with high Si-N bond densities have very low diffusion coefficients. For too high Si-N bond densities the diffusion of the atomic hydrogen in the  $SiN_x:H$  layer will be too slow, resulting in less passivation during short anneals. The similar H-diffusion from the  $SiN_x:H$  layer explains the fact that we find optimum values for bulk passivation for the same Si-N bond densities of  $1.3 \cdot 10^{23} cm^{-3}$  for both  $N_2$  and  $NH_3$  nitrides and regardless of the starting wafer quality. Also the best surface passivation is found for both nitrides around this Si-N bond densities of  $1.3 \cdot 10^{23} cm^{-3}$ .

Consequently, the Si-N bond density found via the FTIR analysis is an important parameter for the solar cell characteristics, related to both the bulk and surface passivation of the SiN<sub>x</sub>:H layers, regardless of the type of SiN<sub>x</sub>:H. This study therefore indicates that FTIR analysis of the SiN<sub>x</sub>:H layers gives us a quick and reliable tool to check the quality and properties of different SiN<sub>x</sub>:H layers. This will allow optimization of SiN<sub>x</sub>:H deposition processes and systems without having to make complete solar cells.

## 5 ACKNOWLEDGEMENTS

This work was partly financially supported by NovemSenter in the DEN program, and the FP6 European Crystal Clear project (EC contract SES6-CT-2003-502583).

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